The asymptotic behaviors for super-Brownian motion with absorption

Yaping Zhu

Peking University (Joint work with Professor Zenghu Li)

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Background

Problem formulation

3 Main results

- Survival probability
- Large deviation
- Absorbed mass

Proofs

- Proof of Theorem 1
- Proof of Theorem 4
- Proof of Theorem 7

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Suppose that ψ is a branching mechanism of the form

$$\psi(\lambda) = -\alpha\lambda + \beta\lambda^2 + \int_{(0,\infty)} (e^{-\lambda x} - 1 + \lambda x) \Pi(dx),$$

where $\alpha = -\psi'(0) \in (0,\infty)$, $\beta \ge 0$ and Π is a σ -finite measure on $(0,\infty)$ satisfying $\int_{(0,\infty)} (x \wedge x^2) \Pi(\mathrm{d}x) < \infty$.

Under \mathbb{P}^x , suppose that $B^{-\rho} = \{B_t^{-\rho} : t \ge 0\}$ is a Brownian motion with drift $-\rho$, starting from x > 0.

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 $\mathcal{M}_{F}(0,\infty)$ $(\mathcal{M}_{F}(\mathbb{R}))$: the space of finite measures on $(0,\infty)$ (\mathbb{R}) .

- $\mathcal{X}^{-\rho} = \{\mathcal{X}_t^{-\rho} : t \ge 0\}$: a superprocess with branching mechanism ψ and spatial motion $\{B_t^{-\rho} : t \ge 0\}$.
- $X^{-\rho} = \{X_t^{-\rho} : t \ge 0\}$: a superprocess with branching mechanism ψ and spatial motion $\{B_{t\wedge\tau_0}^{-\rho} : t \ge 0\}$.

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For all $f \in C_b^+(0,\infty)$, $\mu \in \mathcal{M}_F(0,\infty)$,

$$\mathbb{E}_{\mu}\left(\mathrm{e}^{-\left\langle f,X_{t}^{-
ho}
ight
angle}
ight)=\mathrm{e}^{-\left\langle v_{f}\left(\cdot,t
ight),\mu
ight
angle},\quad t\geq0,$$

where $v_f(x, t)$ is the unique positive solution to the equation

$$\mathbf{v}_f(\mathbf{x},t) = \mathbb{E}^{\mathbf{x}}[f(B_t^{-
ho}); au_0 > t] - \mathbb{E}^{\mathbf{x}}\left[\int_0^t \psi(\mathbf{v}_f(B_s^{-
ho},t-s))\mathrm{d}s; au_0 > t
ight].$$

For $f\in C^+_0(0,\infty)$, the above equation is equivalent to

$$\begin{cases} \frac{\partial}{\partial t} v_f(x,t) = \frac{1}{2} \frac{\partial^2}{\partial x^2} v_f(x,t) - \rho \frac{\partial}{\partial x} v_f(x,t) - \psi(v_f(x,t)), \\ v_f(0+,t) = 0, \quad v_f(x,0) = f(x), \quad t \ge 0, \quad x > 0. \end{cases}$$
(1)

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- Watanabe (1968): existence.
- Li and Shiga (1995): immigration structure, the range of immigration process.
- Li (1999): conditional law.
- Kyprianou, Murillo-Salas and Pérez (2012): the right-most point in the support, absorbed mass.

- Engländer (2004): large deviation.
- Kyprianou, Liu, Murillo-Salas and Ren (2012): additive martingale and derivative martingale.
- Kyprianou and Murillo-Salas (2013): L^p-convergence of additive martingale.
- Ren, Song and Zhang (2021): supermum of the support, extremal process.

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Define $R_t := \inf\{y > 0 : X_t^{-\rho}(y,\infty) = 0\}$ and denote the extinction time of $X^{-\rho}$ by

$$\zeta := \inf\{t > 0 : ||X_t^{-\rho}|| = 0\}.$$

(1) For all $\rho \ge \sqrt{2\alpha}$ and x > 0, $\mathbb{P}_{\delta_x}(\zeta < \infty) = 1$. (2) Assume that $-\infty < \rho < \sqrt{2\alpha}$. Then for all x > 0,

$$\lim_{t\to\infty}\frac{R_t}{t}=\sqrt{2\alpha}-\rho,\quad\text{on }\{\zeta=\infty\}\quad\mathbb{P}_{\delta_x}-\text{a.s.}$$

Problem 1: When $\rho \ge \sqrt{2\alpha}$, what's the decay rate of $\mathbb{P}_{\delta_{\chi}}(\zeta > t)$? Consider the behavior of $||X_t^{-\rho}||$ conditioned on $\zeta > t$.

Problem 2: When $\rho < \sqrt{2\alpha}$, what's the decay rate of $\mathbb{P}_{\delta_x}(R_t > \gamma t)$ for $\gamma > \sqrt{2\alpha} - \rho$? Consider the limiting behavior of $X_t^{-\rho}(\gamma t, \infty)$ conditioned on $R_t > \gamma t$.

Remark: Analogue results for branching Brownian motion (BBM) with absorption can be found in Harris et al. (2006, 2007).

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Problem formulation

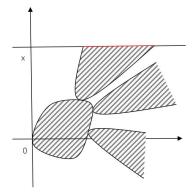


Figure 1. $||X_{D_x}^{-\rho}||$ denote the absorbed mass at x.

When $\rho = -\sqrt{2\alpha}$, $\mathbb{P}_{\delta_0}(||X_{D_x}^{-\rho}|| > t) \sim \frac{\sqrt{2\alpha}xe^{\sqrt{2\alpha}x}}{t(\log t)^2}, \quad t \to \infty.$

Furthermore, when $\rho < -\sqrt{2\alpha}$,

$$\mathbb{P}_{\delta_0}ig(||X_{D_x}^{-
ho}||>tig)\sim Ct^{-d},\quad t o\infty,$$

where $C = C(x, \rho, \psi)$, $d = \overline{\lambda}_{\rho}/\lambda_{\rho}$ and $\lambda_{\rho} \leq \overline{\lambda}_{\rho}$ are the two roots of the quadratic equation $\lambda^2 + 2\rho\lambda + 2\alpha = 0$.

Problem 3: For any $\rho > -\sqrt{2\alpha}$, what's the tail behavior of the absorbed mass?

Remark: Maillard (2013), Lalley and Zheng (2015), Berestycki et al. (2017), Corre (2018): the number of absorbed particles in BBM with absorption.

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Assume that the branching mechanism $\boldsymbol{\psi}$ satisfies

$$\int^{\infty} \frac{1}{\sqrt{\int_{\lambda^*}^r \psi(u) \mathrm{d}u}} \mathrm{d}r < \infty, \tag{2}$$

where λ^* is the largest root of the equation $\psi(\lambda) = 0$.

Remark 1

Let \mathcal{R} be the smallest closed set such that $\mathrm{supp}\mathcal{X}_t^{-\rho}\subseteq \mathcal{R}$, $t\geq 0$. It follows from Sheu (1997) that

$$\mathbb{P}_{\mu}\left(\mathcal{R} ext{ is compact}
ight) = \mathrm{e}^{-\lambda^{st}||\mu||},$$

and (2) implies Grey's condition

$$\int^{\infty} \frac{1}{\psi(\lambda)} \mathrm{d}\lambda < \infty.$$

Furthermore, we also assume the following integrability condition:

$$\int_{0+}\lambda^{-1}\varphi(\lambda)\mathrm{d}\lambda<\infty,$$

where $\varphi(\lambda) := \lambda^{-1} \psi(\lambda) + \alpha$.

Remark 2

According to Kyprianou, Liu, Murillo-Salas and Ren (2012),

$$\int_{0+}\lambda^{-1}\varphi(\lambda)\mathrm{d}\lambda<\infty\Leftrightarrow\int_{[1,\infty)}\lambda(\log\lambda)\mathsf{\Pi}(\mathrm{d}\lambda)<\infty.$$

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Theorem 1 (Survival probability)

For any $\rho > \sqrt{2\alpha}$ and x > 0, we have

$$\lim_{t\to\infty}\mathbb{P}_{\delta_{x}}(\zeta>t)\frac{\sqrt{2\pi t^{3}}}{xe^{\rho x}}e^{(\frac{1}{2}\rho^{2}-\alpha)t}=C_{1},$$
(3)

where C_1 is a positive constant independent of x.

Theorem 2 (Yaglom limit theorem)

For any $\rho > \sqrt{2\alpha}$ and x > 0, there exists a probability distribution $\mathbb Q$ defined on $(0,\infty)$ such that

$$\lim_{t\to\infty}\mathbb{P}_{\delta_{\mathsf{x}}}(||X_t^{-\rho}||\in\cdot|\zeta>t)=\mathbb{Q}(\cdot),$$

and \mathbb{Q} has mean $\frac{2}{C_1 \rho^2}$.

Theorem 3 (Weak convergence)

For any x>0 and $r\geq 0$, there exists a probability distribution \mathbf{Q}_r defined on $\mathcal{M}_F(0,\infty)$ such that

$$\mathbb{P}_{\delta_{x}}(X_{t}^{-\rho} \in \cdot | \zeta > t + r) \xrightarrow{w} \mathbf{Q}_{r}(\cdot).$$

(i) When r = 0, for any $f \in C_b^+(0, \infty)$,

$$\lim_{t\to\infty}\mathbb{P}_{\delta_x}\Big(\mathrm{e}^{-\langle f,X_t^{-\rho}\rangle}\Big|\zeta>t\Big)=1-C_1^{-1}\int_0^\infty 2zf(z)g_f(z)\mathrm{e}^{-\rho z}\mathrm{d} z$$

where C_1 is defined by (3) and for any z > 0,

$$g_f(z) := \mathbb{E}_B^z \left[\exp \left\{ - \int_0^\infty \varphi(v_f(Z(r), r)) \mathrm{d}r
ight\}
ight] \in (0, 1],$$

here $\{Z(t)\}_{t\geq 0}$ is a Bessel-3 process starting from z.

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Theorem 3 (Weak convergence)

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(i) When r = 0, for any $f \in C_b^+(0,\infty)$,

$$\lim_{t\to\infty}\mathbb{P}_{\delta_{X}}\left(\mathrm{e}^{-\langle f,X_{t}^{-\rho}\rangle}\Big|\zeta>t\right)=1-C_{1}^{-1}\int_{0}^{\infty}2zf(z)g_{f}(z)\mathrm{e}^{-\rho z}\mathrm{d}z,$$

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Theorem 3 (Weak convergence)

(ii) For any r>0 and $f\in C_b^+(0,\infty)$, we have

$$\begin{split} &\lim_{t\to\infty} \mathbb{P}_{\delta_{x}}\left(\mathrm{e}^{-\langle f,X_{t}^{-\rho}\rangle}\Big|\zeta>t+r\right)\\ &=\mathrm{e}^{(\frac{1}{2}\rho^{2}-\alpha)r}C_{1}^{-1}\int_{0}^{\infty}2z\mathrm{e}^{-\rho z}\left[\left(f(z)+u(z,r)\right)g_{f}(z,r)-f(z)g_{f}(z)\right]\mathrm{d}z, \end{split}$$

where

$$g_f(z,r) := \mathbb{E}_B^z \left[\exp\left\{ -\int_0^\infty \varphi(v_{f+u_r}(Z(s),s)) \mathrm{d}s
ight\}
ight], \quad z > 0.$$

and $u_t(x) := u(x,t) = -\log \mathbb{P}_{\delta_x}(\zeta \leq t).$

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Theorem 4 (Large deviation)

For any $\gamma > \sqrt{2\alpha} - \rho$ and x > 0, we have

$$\lim_{t\to\infty}\mathbb{P}_{\delta_{\mathrm{x}}}(\mathsf{R}_t>\gamma t)\frac{\sqrt{2\pi t}}{1-\mathrm{e}^{-2\gamma\mathrm{x}}}\mathrm{e}^{(\frac{1}{2}(\gamma+\rho)^2-\alpha)t}\mathrm{e}^{-(\gamma+\rho)\mathrm{x}}=\mathsf{C}_2,$$

where C_2 is a positive constant independent of x.

Theorem 5 (Yaglom limit theorem)

For any $\gamma > \sqrt{2\alpha} - \rho$ and x > 0, there exists a probability distribution Π such that

$$\lim_{t\to\infty}\mathbb{P}_{\delta_x}(X_t^{-\rho}(\gamma t,\infty)\in\cdot|R_t>\gamma t)=\mathbf{\Pi}(\cdot).$$

Moreover, Π has mean $\frac{1}{C_2(\gamma+\rho)}$.

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Define
$$\Theta_{\gamma t} X_t^{-\rho}(\cdot) = X_t^{-\rho}(\cdot + \gamma t)$$
 and $Y_t^{-\rho} := \Theta_{\gamma t} X_t^{-\rho}$.

Theorem 6 (Weak convergence)

Assume $\gamma > \sqrt{2\alpha} - \rho$. For any x > 0 and $s \ge 0$, there exists a probability distribution Π_s on $\mathcal{M}_F(\mathbb{R})$ such that

$$\mathbb{P}_{\delta_{x}}(Y_{t}^{-\rho} \in \cdot | R_{t+s} > \gamma t) \xrightarrow[t \to \infty]{w} \Pi_{s}(\cdot).$$

(i) When s = 0, for any $f \in C_b^+(0, \infty)$,

$$\lim_{t\to\infty}\mathbb{P}_{\delta_{\mathbf{x}}}\left(\mathrm{e}^{-\langle f,Y_t^{-\rho}\rangle}\Big|R_t>\gamma t\right)=1-C_2^{-1}\int_0^\infty f(z)h_f(z)\mathrm{e}^{-(\rho+\gamma)z}\mathrm{d}z,$$

where

$$h_f(z) := \mathbb{E}\left(\mathrm{e}^{-\int_0^\infty \varphi(v_f(W_2(r)+z-\gamma r,r))\mathrm{d}r}
ight), \quad z>0,$$

and $\{W_2(t)\}_{t\geq 0}$ is a standard Brownian motion.

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(i) When s = 0, for any $f \in C_b^+(0,\infty)$,

$$\lim_{t\to\infty}\mathbb{P}_{\delta_{x}}\left(\mathrm{e}^{-\langle f,Y_{t}^{-\rho}\rangle}\Big|R_{t}>\gamma t\right)=1-C_{2}^{-1}\int_{0}^{\infty}f(z)h_{f}(z)\mathrm{e}^{-(\rho+\gamma)z}\mathrm{d}z,$$

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$$h_f(z) := \mathbb{E}\left(\mathrm{e}^{-\int_0^\infty \varphi(v_f(W_2(r)+z-\gamma r,r))\mathrm{d}r}\right), \quad z>0,$$

and $\{W_2(t)\}_{t\geq 0}$ is a standard Brownian motion.

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Theorem 6 (Weak convergence)

(ii) For any s > 0 and $f \in C_b^+(0,\infty)$, we have

$$\begin{split} &\lim_{t\to\infty} \mathbb{P}_{\delta_x} \left(\mathrm{e}^{-\langle f, Y_t^{-\rho} \rangle} \Big| R_{t+s} > \gamma t \right) \\ &= C_2^{-1} \mathrm{e}^{\frac{1}{2}(\gamma+\rho)^2 s - \alpha s} \int_0^\infty \left[\left(f(z) + u^0(z,s) \right) h_f(z,s) - f(z) h_f(z) \right] \mathrm{e}^{-(\rho+\gamma)z} \mathrm{d}z, \end{split}$$

and for any z > 0,

$$h_f(z,s) = \mathbb{E}\left[e^{-\int_0^\infty \varphi\left(v_{f+u_s^0}(W_2(r)+z-\gamma r,r)\right)dr}\right] \in (0,1],$$

where $u^0_t(x):=u^0(x,t)=-\log \mathbb{P}_{\delta_x}(\mathcal{X}^{ho}_t(0,\infty)=0).$

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Define

$$u_s(x) := -\log \mathbb{E}_{\delta_0}\left(\mathrm{e}^{-s||X_{D_x}^{-
ho}||}
ight), \quad x > 0, s \in \mathbb{R},$$

and the singular point

$$s_0 := \inf\{s < 0 : u_s > -\infty\}.$$

Theorem 7 (Tail distribution of the absorbed mass)

When $ho \geq \sqrt{2 lpha}$, we have $-\infty < {\it s}_0 < {\it 0}$ and

$$\mathbb{P}_{\delta_0}\left(||X_{D_x}^{-
ho}||>t
ight)\sim rac{u_{s_0}'(x)\mathrm{e}^{-u_{s_0}(x)}}{\mathrm{e}^{-s_0t}\sqrt{\pi t\psi(s_0)}},\quad t
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Theorem 1 (Survival probability)

For any $\rho > \sqrt{2\alpha}$ and x > 0, we have

$$\lim_{t\to\infty}\mathbb{P}_{\delta_x}(\zeta>t)\frac{\sqrt{2\pi t^3}}{xe^{\rho x}}e^{(\frac{1}{2}\rho^2-\alpha)t}=C_1,$$

where C_1 is a positive constant independent of x.

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Step 1. According to the Laplace transform of $X^{-\rho}$,

$$\mathbb{P}_{\delta_{\mathsf{X}}}(\zeta > t) = 1 - \lim_{ heta o \infty} \lim_{n o \infty} e^{-u_n^{ heta}(\mathsf{X},t)},$$

where $u_n^{\theta}(x, t)$ is the solution to (1) with boundary condition $u_n^{\theta}(x, 0) = \theta f_n(x)$ and $f_n \in C_0^+(0, \infty)$ satisfying $f_n \uparrow 1_{(0,\infty)}$ as $n \uparrow \infty$.

It follows from monotone convergence theorem that

$$\mathbb{P}_{\delta_x}(\zeta > t) = 1 - \mathrm{e}^{-u(x,t)}.$$

Step 2. Since $\{||\mathcal{X}_t^{-\rho}|| : t \ge 0\}$ is a continuous state branching process (CSBP), naturally we can use the cumulant semigroup of CSBP to give an upper bound for u(x, t).

Lemma

For any x > 0,

$$u(x,t) \leq \hat{v}_t, \quad t > 0,$$

where $\hat{v}_t := \uparrow \lim_{\theta \to \infty} v_t(\theta)$ and $t \mapsto v_t(\theta)$ is the unique positive solution of

$$\mathbf{v}_t(heta) = heta - \int_0^t \psi(\mathbf{v}_s(heta)) \mathrm{d}s, \quad t \ge 0.$$

Moreover, for any t > 0, $\hat{v}_t < \infty$ if and only if Grey's condition holds.

Step 3. It follows from Feynman-Kac formula that

$$\begin{split} u(x,t) \frac{\sqrt{2\pi t^3}}{xe^{\rho x}} \mathrm{e}^{(\frac{1}{2}\rho^2 - \alpha)(t-1)} \\ &= \int_{(0,\infty)} \frac{\sqrt{t}}{\sqrt{t-1}} \frac{t}{x} (1 - e^{-\frac{2xz}{t-1}}) e^{-\rho z} u(z,1) \\ &\times \mathbb{E} \left[e^{-L_{z,x}(1,t)} | \tau_0^{t-1}(z,x) > t-1 \right] e^{-\frac{(x-z)^2}{2(t-1)}} \mathrm{d}z \\ &\to \int_{(0,\infty)} 2z e^{-\rho z} u(z,1) \mathbb{E}_B^z (e^{-\int_1^\infty \varphi(u(Z(r-1),r)) \mathrm{d}r}) \mathrm{d}z \end{split}$$

where $L_{z,x}(1,t) = \int_1^t \varphi(u(B_{z,x}^{t-1}(r-1),r)) dr.$

Proposition

For each z > 0, we have

$$\lim_{t\to\infty}\mathbb{E}\big[e^{-\mathcal{L}_{z,x}(1,t)}|\tau_0^{t-1}(z,x)>t-1\big]=\mathbb{E}_B^z(e^{-\int_1^\infty \varphi(u(Z(r-1),r))\mathrm{d} r})\in(0,1],$$

where $\{Z(s) : s \ge 0\}$ is a Bessel-3 process starting from z under \mathbb{P}_B^z .

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Theorem 4 (Large deviation)

For any $\gamma > \sqrt{2\alpha} - \rho$ and x > 0, we have

$$\lim_{t\to\infty}\mathbb{P}_{\delta_{\mathbf{x}}}(R_t>\gamma t)\frac{\sqrt{2\pi t}}{1-\mathrm{e}^{-2\gamma\mathbf{x}}}\mathrm{e}^{(\frac{1}{2}(\gamma+\rho)^2-\alpha)t}\mathrm{e}^{-(\gamma+\rho)\mathbf{x}}=C_2,$$

where C_2 is a positive constant independent of x.

Step 1.
$$\mathbb{P}_{\delta_{X}}(R_{t} > \gamma t) = 1 - \exp\{-v^{\gamma t}(x, t)\}.$$

Step 2. For any x, t > 0, $v^{\gamma t}(x, t) \le \hat{v}_t < \infty$.

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Step 3. It follows from Feynman-Kac formula that

$$v^{\gamma t}(x,t) = \frac{\mathrm{e}^{\rho x - (\rho^2/2 - \alpha)(t-1)}}{\sqrt{2\pi(t-1)}} \int_{(0,\infty)} v^{\gamma t}(z,1) \mathrm{e}^{-\frac{(x-z)^2}{2(t-1)}} \mathrm{e}^{-\rho z} \\ \times \mathbb{E} \big[\mathrm{e}^{-J_{z,x}(0,t-1)}; \tau_0^{t-1}(z,x) > t-1 \big] \mathrm{d}z,$$

where $J_{z,x}(0, t-1) = \int_0^{t-1} \varphi(v^{\gamma t}(B_{z,x}^{t-1}(r), r+1)) dr$.

Since the effect of the absorption vanishes as the start position tends to infinity, thus

$$\mathbf{v}^{\gamma t}(z+\gamma t,1)
ightarrow u^0(z,1):= -\log \mathbb{P}_{\delta_x}ig(\mathrm{e}^{-\mathcal{X}_t(0,\infty)}ig), \quad t
ightarrow \infty.$$

Proposition

For any
$$\gamma > \sqrt{2\alpha} - \rho$$
, $x > 0$ and $z \in \mathbb{R}$, we have

 $\lim_{t\to\infty}\mathbb{E}\big[\mathrm{e}^{-J_{z+\gamma t,x}(0,t-1)};\tau_0^{t-1}(z+\gamma t,x)>t-1\big]=(1-\mathrm{e}^{-2\gamma x})h(z),$

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where $h : \mathbb{R} \to (0, 1]$ is a strictly positive function.

Theorem 7 (Tail distribution of the absorbed mass)

When $ho \geq \sqrt{2lpha}$, we have that $-\infty < s_0 < 0$ and

$$\mathbb{P}_{\delta_0}\left(||X_{D_x}^{-
ho}||>t
ight)\sim rac{u_{s_0}'(x)\mathrm{e}^{-u_{s_0}(x)}}{\mathrm{e}^{-s_0t}\sqrt{\pi t\psi(s_0)}},\quad t
ightarrow\infty.$$

Recall that

$$u_{s}(x):=-\log \mathbb{E}_{\delta_{0}}\Big(\mathrm{e}^{-s||X_{D_{x}}^{-
ho}||}\Big), \quad x>0,s\in \mathbb{R}.$$

According to Kyprianou et al. (2012a), when $\rho \ge \sqrt{2\alpha}$, $\{||X_{D_x}^{-\rho}|| : x \ge 0\}$ is a subcritical CSBP and $u_s(x)$ satisfies the following equation

$$\frac{1}{2}u_{s}''(x) + \rho u_{s}'(x) - \psi(u_{s}(x)) = 0, \quad x > 0,$$

with boundary conditions $u_s(0) = s$ and $u_s(\infty) = 0$.

Based on classical results on some complex differential equations, we analyze the behavior of u_s near its singular point s_0 .

Lemma 1

When
$$\rho \geq \sqrt{2\alpha}$$
, we have $-\infty < s_0 < 0$, $u_{s_0} > -\infty$, and $u_{s_0}'(0) = 0$.

Lemma 2

When $\rho \ge \sqrt{2\alpha}$, for x > 0 fixed, there exists $r_x > 0$ such that $u_s(x)$ is analytic on $V = D(s_0, r_x) \setminus (-\infty, s_0]$ and

$$\partial_s u_s(x) \sim rac{u_{s_0}'(x)}{2\sqrt{\psi(s_0)(s-s_0)}}, \quad s
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References

- Berestycki, J.; Brunet, Éric; Harris, S. C.; Miłoś, P. (2017): Branching Brownian motion with absorption and the all-time minimum of branching Brownian motion with drift. *J. Funct. Anal.* **273**, 2107-2143.
- Harris, J. W.; Harris, S. C.; Kyprianou, A. E. (2006): Further probabilistic analysis of the Fisher-Kolmogorov-Petrovskii-Piscounov equation: one sided travelling-waves. *Ann. Inst. H. Poincaré Probab. Statist.* **42**, 125-145.
- Harris, J. W.; Harris, S. C. (2007): Survival probabilities for branching Brownian motion with absorption. *Electron. Comm. Probab.* **12**, 81-92.



Kyprianou, A. E.; Murillo-Salas, A.; Pérez, J. L. (2012): An application of the backbone decomposition to supercritical super-Brownian motion with a barrier. *J. Appl. Probab.* **3**, 671-684.



- Li, Z. H.; Zhu, Y. P. (2022): Survival probability for super-Brownian motion with absorption. *Statist. Prob. Lett.* **186**, 9 pp.
- Zhu, Y. P. (2023): A large deviation theorem for a supercritical super-Brownian motion with absorption. J. Appl. Probab. 1-26.

Thanks

Email: zhuyp@pku.edu.cn

